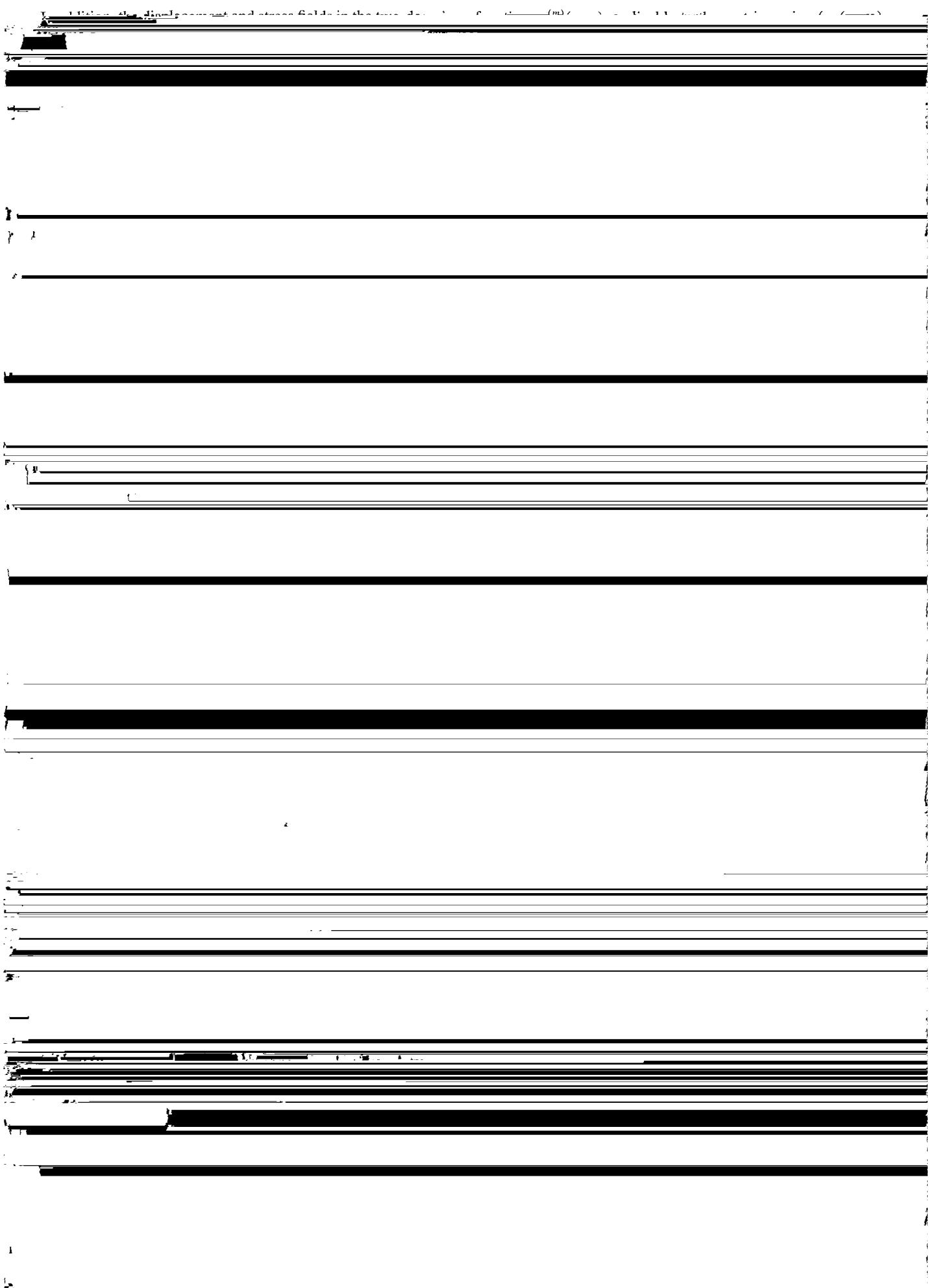


In this paper, we restrict attention to matrix cracking which accompanies the complete fracture of an isolated fiber. It is assumed that this fracture configuration results in a defect which can be modeled as a penny-shaped crack located in a single

$$\sigma_{\theta\theta}^{(m)}(r, z) = -\frac{A_m}{r^2} + 2B_m \quad (1b)$$



$$\eta R_3(s, t) = \sinh(st)(a_7a_{15} - a_{10}a_{12}) - (a_{15}I_0 + a_{10}I_1)\{\sinh(st)(2\nu_f - 1) + st \cosh(st)\} \quad (51)$$

$$\eta R_4(s, t) = \sinh(st)(a_7a_{16} - a_{11}a_{12}) - (a_{16}I_0 + a_{11}I_1)\{\sinh(st)(2\nu_f - 1) + st \cosh(st)\}. \quad (52)$$

Considering the integrals of Eqs. (32) to (35) containing $F_1(u)$, the integral expressions (36) to (39) and the representation (43) it can be shown that

$$-\frac{2}{\pi} \int_0^\infty \{2(1 - \nu_f)f_1 + uf_2\} \frac{F_1(u)}{u} J_0(ua) du$$

$$= -\frac{2}{\pi} \int_0^a [(3 - 2\nu_f) \sinh(su) K_0(sa)$$

$$-saK_1(sa) \sinh (su) + su \cosh (su) K_0(sa)]m(u)du;$$

$$+ sI_1(sa) \left[\int_a^b uh(u)e^{-su}du + \int_b^\infty ug(u)e^{-su}du \right];$$

$$0 < s < \infty \quad (58)$$

$$-\frac{2}{\pi} a \mu_m \int_0^\infty \left[(-f_3 + u f_4) J_0(ua) \right.$$

$$+ \left\{ (1 - 2\nu_m) f_3 - u f_4 \right\} \frac{J_1(ua)}{ua} \Bigg] F_2(u) du$$

$$= \mu_m \left[\int_a^b h(t) e^{-st} \right] \left\{ sa \left\{ \frac{1}{2} I_0(sa) + saI_1(sa) - stI_0(sa) \right\} \right.$$

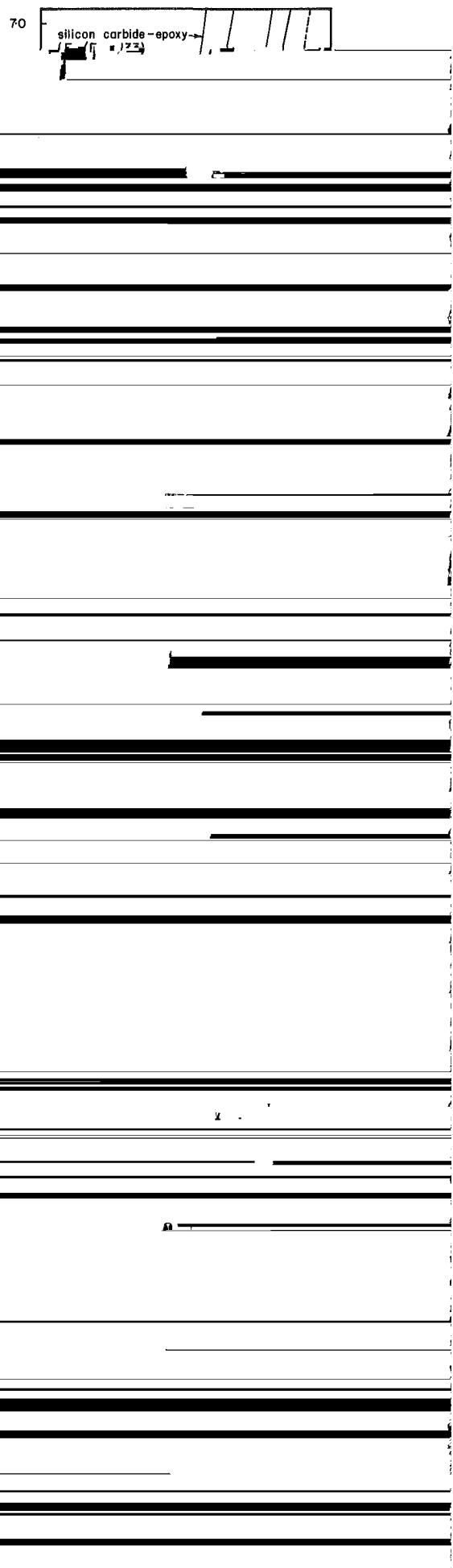
$$\frac{M(t)}{t} = \frac{\pi \mu_m m(t)}{\Gamma_1^2(2\nu_m)} - \frac{s^2 a}{\Gamma_1^2(2\nu_m) + s^2 I_1(sa)}$$

and the kernel functions S_1 and S_2 are given by

$$S_1(u, t) = \int_0^\infty e^{-su} [\{ (3 - 2\nu_m) I_0(sa) + sa I_1(sa) \\ - s^2 u I_1(sa) \} R_4(s, t)] ds. \quad (68)$$

$$S_2(u, t) = - \int_0^\infty [\{ (3 - 2\nu_r) \sinh(su) K_0(sa) \\ - s^2 u K_1(sa) \} R_4(s, t)] ds.$$

and subjected to a uniform strain ϵ_0 (or uniform stress $E_m \epsilon_0$)
(Sneddon, 1946), i.e.,



$$a_1 = s^2 a^2 (I_1^2 - I_0^2) + 2(1 - \nu_f) I_1^2 \quad (\text{A7})$$

$$a_2 = sa\Gamma\{K_0 I_1 + K_1 I_0\} \quad (\text{A8})$$

$$a_3 = -\Gamma[K_1 I_1 \{2 - 2\nu_m + s^2 a^2\} + 2(1 - \nu_m)saI_0K_1] \quad (\text{A9})$$

$$a_4 = sa(I_1^2 - I_0^2) + 4(1 - \nu_f)I_0 I_1 \quad (\text{A10})$$

$$X_1(s) = [(3 - 2\nu_m)I_0(sa)$$

$$+ saI_1(sa)] \left[\int_a^b h(u)e^{-su} du + \int_b^\infty g(u)e^{-su} du \right]$$

$$- sI_0(sa) \left[\int_a^b uh(u)e^{-su} du + \int_b^\infty ug(u)e^{-su} du \right]$$

$$a_6 = -\{sa(I_1K_1 + I_0K_0) + 4(1 - \nu_m)I_0K_1\} \quad (\text{A12})$$

$$+ su \cosh (su)K_0(sa)]m(u)du \quad (\text{A24})$$

$$a_7 = -\frac{1}{8\pi I} \left\{ saK_1(a_2a_4 - a_5a_1) + K_0(a_3a_4 - a_6a_1) \right\} \quad (\text{A13}) \quad X_2(s) = -\left\{ \frac{sa}{8\pi I} [I_0(sa) + I_2(sa)] + 2\nu_m I_1(sa) \right\}$$

$$a_8 = 1 + \frac{I_1(a_2saK_1 + a_3K_0)}{(a_2a_6 - a_3a_5)} \quad (\text{A14})$$

$$\times \left\{ \int_a^b h(u)e^{-su} du + \int_b^\infty g(u)e^{-su} du \right\}$$

$$a_9 = \frac{I_0(sa a_2 K_1 + a_3 K_0)}{(a_2 a_6 - a_3 a_5)} \quad (\text{A15})$$

$$+ sI_1(sa) \left[\int_a^b uh(u)e^{-su} du + \int_b^\infty ug(u)e^{-su} du \right]$$

The expressions for $P_i(s, t)$ ($i = 1, 2, 3, 4$) occurring in the kernel functions S_1 and S_2 defined by Eqs. (65) and (66) are given by

$$\zeta P_1(s, t) = -[\{(a_{12}I_0 + a_7I_1)a_3I_1 + (a_8I_1 + a_{13}I_0)$$

where the unknown functions $X(t)$ are defined as

$$X(t) = \begin{cases} M(t); & 0 \leq t < a \\ H(t); & a \leq t < b \end{cases} \quad (B2)$$

$$- (a_2a_4 - a_1a_5)(a_{13}I_0 + a_8I_1) \} \\ \times \{ 2\nu_m L_1(s, t) - sL_2(s, t) \}] \quad (A28)$$

$$\zeta P_2(s, t) = -[\{(a_{12}I_0 + a_7I_1)a_3I_0 + (a_9I_1 + a_{14}I_0)$$

$$- (a_2a_4 - a_1a_5)(a_{14}I_0 + a_9I_1) \}$$

The right-hand side of the prescribed function $\mathbf{B}(t)$ in (B1) takes the form

$$\mathbf{B}(t) = \begin{cases} t\Gamma\Omega; & 0 \leq t < a \\ (b^2 - t^2)^{1/2}; & a \leq t < b \\ 0; & b \leq t < \infty \end{cases} \quad (B3)$$

$$- (a_2a_4 - a_1a_5)(a_{14}I_0 + a_9I_1) \}$$

where $N = N_1 + N_2 + N_3$. Using the above procedures and representations, the discretized form of (B1) can be written as

$$\mathbf{A}_x \mathbf{x}_x = \mathbf{R}_x \quad (B11)$$

$$\mathbf{A}_{ij} = \delta_{ij} + K(t_i, t_j) \Delta t, \quad (B14)$$

where δ_{ij} is the Kronecker delta and Δt is given as h_x , h_y or